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# Matrix Isolation and Theoretical Study on the Photolysis of CH<sub>2</sub>CICOCI

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Authors' contributions

This work was carried out in collaboration between both authors. Author NT performed the experimental and computational studies, wrote the protocol and wrote the first draft of the manuscript. Author MN performed the experimental study. Both authors read and approved the final manuscript.

Original Research Article

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# ABSTRACT

UV light photolysis of chloroacetyl chloride (CH<sub>2</sub>CICOCI) has been investigated using infrared spectroscopy in cryogenic Ar, Kr, Xe and O<sub>2</sub> matrices. In all matrices, CHCI=C=O, HCI, CH<sub>2</sub>Cl<sub>2</sub> and CO were produced, with the relative yields depending on the used matrix gas. The relative yield of CH<sub>2</sub>Cl<sub>2</sub> formation to CHCI=C=O formation was greatly enhanced in Xe because of the external heavy atom effect, indicative of the formation of CH<sub>2</sub>Cl<sub>2</sub> via the C–C or C–CI bond cleavage in the triplet state. In CH<sub>2</sub>CICOCI/Ar, photoisomerization from *anti*- to *gauche*-CH<sub>2</sub>CICOCI was also observed at the early stage of the irradiation. The results show that in Ar, the reactions predominantly proceed through concerted mechanisms.

Keywords: Chloroacetyl chloride; photolysis; cryogenic matrix; Ketene.

# **1. INTRODUCTION**

Chlorinated acetyl chlorides are produced through the oxidation of chlorinated ethenes [1,2]. The photodissociation of chlorinated acetyl chlorides is an environmentally important

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process. The potential release of CI by photochemical reaction in the atmosphere enhances the CI cycle in ozone depletion. In the oxidation of chlorinated ethenes initiated by CI, relatively high yields of chlorinated acetyl chloride were reported by Hasson and Smith [3]. In Ar matrix, hydroxyl radical reacts with tetrachloroethene to yield trichloroacetyl chloride [4].

Chloroacetyl chloride can exist as an *anti* or *gauche* conformer depending on the difference in the relative orientation of the two C–Cl bonds [5]. Infrared (IR) spectra of chloroacetyl chloride in gas, liquid and solid phases, and matrix-isolated were measured by many groups [6-10]. Allen and Russell studied the pyrolysis of chloroacetyl chloride using IR laser, in which 1,2-HCl elimination occurred to produce chloroketene [11]. In Xe matrix, Davidovics et al. obtained the photolysis product of chloroketene [12]. Previously, we studied the photolysis of CCl<sub>3</sub>COCl in Ar and O<sub>2</sub> matrices, in which C–C bond cleavage was found to be the major reaction path [13]. In contrast, acetyl chloride (CH<sub>3</sub>COCl) underwent the fourcenter elimination to yield the HCl-ketene complex in the S<sub>0</sub> state after internal conversion from the S<sub>1</sub> state [14,15]. For the photolysis of CF<sub>3</sub>COCl, CF<sub>3</sub>Cl and CO were produced via the radical mechanism [16]. The fact that the substitution of a halogen atom could change the reaction mechanism is of great interest.

In the present study, we investigated the UV light photolysis of CH<sub>2</sub>CICOCI in cryogenic Ar, Kr, Xe and O<sub>2</sub> matrices with the aid of theoretical calculations using the B3LYP and MP2 methods. The computational results were used to interpret the IR spectra. Moreover, the reaction mechanism was discussed on the basis of the differences observed in the four matrices.

## 2. EXPERIMENTAL DETAILS

Light irradiation was performed using a low-pressure mercury short arc lamp (HAMAMATSU L937-04,  $\lambda > 253.7$  nm). IR spectra were measured in the range of 4000–700 cm<sup>-1</sup> with a resolution of 1.0 cm<sup>-1</sup> using a Fourier transform IR spectrometer (SHIMADZU 8300A) coupled with a liquid-nitrogen-cooled MCT detector. Each spectrum was obtained by acquiring 128 scans. A closed-cycle helium cryostat (Iwatani M310/CW303) was used to control the temperature of the matrix.

Argon (Nippon Sanso, 99.9999%), krypton (Taiyo Sanso), xenon (Nippon Sanso) and  $O_2$  (Okaya Sanso) were used without further purification. Chloroacetyl chloride (Wako Pure Chemicals) was used after freeze–pump–thaw cycle at 77 K. Dichloromethane (Wako Pure Chemicals) was used as an authentic sample for product identification. The sample was diluted with the matrix gas to approximately 1/1000 (0.2 Torr sample and 200 Torr matrix gas) and was deposited on a Csl window at 6 K.

For product identification and energetic consideration, molecular orbital calculation was performed. Geometry optimizations were performed using the second-order Møller–Plesset theory (MP2) with the 6-311++G(3df,3pd) basis set and using the Becke's three-parameter hybrid density functional [17] in combination with the Lee–Yang–Parr correlation functional (B3LYP) [18] with the aug-cc-pV(T+d)Z basis set. Harmonic vibrational frequency calculation was performed to confirm the predicted structures as local minima and to elucidate the zero-point vibrational energy corrections. The vertical transition energy was calculated at the SAC-CI/D95+(d,p) level based on the structures optimized at the CCSD/D95+(d,p) level. All calculations were performed using Gaussian 09 [19].

#### 3. RESULTS AND DISCUSSION

#### 3.1 CH<sub>2</sub>CICOCI/Ar

A mixture of CH<sub>2</sub>CICOCI/Ar was deposited on a CsI window with a ratio of CH<sub>2</sub>CICOCI/Ar = 1/1000. In the IR spectrum obtained after deposition, a strong band was observed at 1829 cm<sup>-1</sup>, which was attributed to the C=O stretching vibration of the anti-CH<sub>2</sub>CICOCI [6,9]. Fig. 1(a) shows the IR difference spectrum obtained after irradiation at a wavelength >253.7 nm of the CH<sub>2</sub>CICOCI/Ar matrix for 210 min. The positive and negative bands indicate the growth and depletion, respectively, during the irradiation period. Table 1 lists the observed wavenumbers of the growth bands. Although the band at 1790 cm<sup>-1</sup>, which was previously assigned to the C=O stretching vibration of gauche-CH2CICOCI, decreased concomitantly with the increase of the shoulder at  $1784 \text{ cm}^{-1}$ , the growth bands at 1140, 1060, 1054, 822, 813 and 810 cm<sup>-1</sup> are in good agreement with those of gauche-CH<sub>2</sub>CICOCI [9]. Based on far-IR spectrum measurements, Durig et al. determined that the enthalpy difference between the two rotamers was 587 cm<sup>-1</sup> [6], at which the anti rotamer was more stable. The calculation at the B3LYP/aug-cc-pV(T+d)Z level indicated that the *anti* rotamer is more stable than the *gauche* rotamer by 599 cm<sup>-1</sup> ( $\Delta G^{\circ}$  at 298 K), indicating that the *gauche/anti* population ratio before UV irradiation is 0.086/1 at 298 K. However, the population ratio was estimated to be 0.17 using the relative integrated absorbance ratio  $(I_{1790}/I_{1829})$  of 0.20 and the calculated IR intensity ratio of 1.15 at the B3LYP/aug-cc-pV(T+d)Z level, which is larger than that expected from the Boltzmann distribution. The frequency calculation with the anharmonic correction at the MP2/6-311++G(3df,3pd) level revealed that the band at 1790 cm<sup>-1</sup> was assigned to the overtone band of the CH<sub>2</sub> rocking vibration of anti-CH<sub>2</sub>CICOCI. The barrier height for the conversion from the anti to gauche rotamer was calculated to be 1098 cm<sup>-1</sup> in the S<sub>0</sub> ground state, indicating that the conversion is not expected to occur at 7 K in the absence of UV irradiation. UV irradiation yielded an increase in the population of the less stable rotamer. Further irradiation enabled to distinguish the product bands due to different growth behaviors. Fig. 1(b) shows the IR difference spectrum obtained by subtracting the spectrum measured after 210 min irradiation from that measured after 630 min irradiation. In addition to the anti-CH<sub>2</sub>CICOCI depletion, the bands due to gauche- $CH_2CICOCI$  decreased. The band at 1784 cm<sup>-1</sup>, which shows depletion in Fig. 1(b), was much broader than that in Fig. 1(a), which enabled us to reconfirm that the band at 1784 cm<sup>-1</sup> corresponded to the C=O stretching vibration band of gauche-CH<sub>2</sub>CICOCI overlapping with the band due to anti-CH<sub>2</sub>CICOCI. Prominent doublet bands at 2153 and 2147 cm<sup>-1</sup> continued to grow at different growth rates during the prolonged irradiation period. Three doublet bands at 1302 and 1297  $\text{cm}^{-1}$ , 1120 and 1113  $\text{cm}^{-1}$ , and 808 and 804  $\text{cm}^{-1}$  showed the same growth behavior as those of the 2153 and 2147 cm<sup>-1</sup> bands which were assigned to the C=C stretching, C-H in-plane bending, C-CI stretching and C=O stretching vibrations of CHCI=C=O, respectively. Davidovics et al. measured the IR bands of CHCI=C=O at 2141.4, 1293, 1107 and 842.2 cm<sup>-1</sup> in Xe matrix [12]. The bands at 2857 and 2809 cm<sup>-1</sup> were assigned to the stretching vibration of HCI and their growth behaviors corresponded to those of the bands at 2147 and 2153 cm<sup>-1</sup> of CHCI=C=O, respectively. The weak bands at 1268 and 749 cm<sup>-1</sup> were assigned to the CH<sub>2</sub> wagging and CCl<sub>2</sub> antisymmetric stretching vibrations of CH<sub>2</sub>Cl<sub>2</sub>, respectively, by comparing with the wavenumbers of authentic CH<sub>2</sub>Cl<sub>2</sub> in Ar. Moreover, the weak bands in the 3000–2700 cm<sup>-1</sup> region, except those of HCl were determined as the combination and overtone bands. In contrast to the photolysis of CCI<sub>3</sub>COCI, the band due to COCI was not observed.

Wavenumber/cm <sup>-1</sup>				species
Ar	Kr	Хе	<b>O</b> <sub>2</sub>	
3134			3132	CHCI=C=O <sup></sup> HCI π complex
3115			3115	CHCI=C=OHCI n complex
2885				HCI
2857	2843	2826	2850	HCI <sup></sup> CHCI=C=O n complex
2816/2809	2810/2804	2808	2805	HCI <sup></sup> CHCI=C=O π complex
	2762			CHCl <sub>2</sub> [20]
			2343	<sup>12</sup> CO <sub>2</sub>
			2276	<sup>13</sup> CO <sub>2</sub>
2153	2149	2146	2157	CHCI=C=O···HCI π complex
2147	2145	2142	2152	CHCI=C=OHCI n complex
2140	2142	2133	2145	CO
			2107	$O_3 v_1 + v_3 [21]$
2101				<sup>13</sup> CO
2090	2097			CHCI=C=O
4704			2037	$CO_3$
1784				gauche-CH <sub>2</sub> CICOCI overlapped with
1202	1202	1202		
1302	1303	1295		
1297	4007	1000	4000	
1268	1267	1263	1268	
1140			1130	
1120		4440	4440	
1113		1112	1112	CHCI=C=O-HCI n complex
1109				CHCI=C=O-HCI n complex
4000			1101	$O_3 v_1 [22]$
1060			1057	
1054			1037	
		054	1030	$O_3 V_3 [22]$ $Y_{0} H^+ [23]$
		904		$Ae_{2}(1)$ [25] CHCl <sub>2</sub> [24]
		892		$CHCl_{2}$ [24]
	853	002		$Kr_{2}H^{+}[23]$
		844		$Xe_{2}H^{+}[23]$
822			820	gauche-CH <sub>2</sub> CICOCI
818			816	gauche-CH <sub>2</sub> CICOCI
813				gauche-CH <sub>2</sub> CICOCI
810			810	gauche-CH <sub>2</sub> CICOCI
808			807	CHCI=C=O <sup></sup> HCI π complex
804			803	CHCI=C=OHCI n complex
748	747	746	750	CH <sub>2</sub> Cl <sub>2</sub>
		731		Xe <sub>2</sub> H <sup>+</sup> [23]

Table 1. FTIR spectra of photochemical reaction products of the CH <sub>2</sub> CICOCI in the A	r,
Kr, Xe and O <sub>2</sub> matrices	



Fig. 1. Infrared difference spectra upon  $\lambda > 253.7$  nm irradiation of the matrix CH<sub>2</sub>CICOCI/Ar = 1/1000. (a) 210–0 min and (b) 630–210 min

Two sets of bands assigned to CHCI=C=O vibrations were measured. They showed different growth behavior accompanying with the bands due to HCI different in wavenumbers, indicating the presence of two different structures of the CHCI=C=O...HCI complex in spite of the matrix site effect. As shown in Fig. 2, two stable structures were calculated for the CHCI=C=O...HCI complexes which the hydrogen atom of HCI points to the carbon atom of CHCI ( $\pi$  complex) or O atom (n complex). Similar structures were proposed for the CH<sub>2</sub>=C=O...HCI complexes [25]. Characteristic HCI stretching and C=O stretching vibrations for the two CHCI=C=O...HCI were calculated to be 2821 and 2177 cm<sup>-1</sup>, respectively, for the  $\pi$  complex, and 2919 and 2172 cm<sup>-1</sup>, respectively, for the n complex at the MP2/6-311++G(3df,3pd) level with anharmonic correction. Two species characterized by the absorption bands at 3134, 2809, 2153, 1302, 1120 and 808 cm<sup>-1</sup>, and 3115, 2857, 2147, 1297, 1113 and 804 cm<sup>-1</sup> were identified with the  $\pi$  and n complexes of CHCI=C=O...HCI, respectively.

## 3.2 CH<sub>2</sub>CICOCI/Kr and CH<sub>2</sub>CICOCI/Xe

Figs. 3(a) and 4(a) show the IR difference spectra obtained upon irradiation of the matrices CH<sub>2</sub>CICOCI/Kr and CH<sub>2</sub>CICOCI/Xe, respectively, at a wavelength >253.7 nm. The strong absorption bands observed at approximately 2145 cm<sup>-1</sup> were assigned to the CO stretching vibrations of CHCI=C=O and CO. Due to the weak intensities of the bands of CHCI=C=O, except for the C=O stretching band, the presence of the two complexes was not evident, although splittings were observed in the 2145 cm<sup>-1</sup> region. As observed in Figs. 2(b) and 3(b), the intensity of the higher wavenumber side of these bands decreased. The product CHCI=C=O was further photodissociated in the Kr and Xe matrices. SAC-CI calculation showed that the S<sub>1</sub> state of CHCI=C=O possesses the mixing characters of  $\pi\sigma_{C-CI}^{*}$  and  $\pi$ Rydberg, -0.59 (HOMO  $\rightarrow$  LUMO) - 0.53 (HOMO  $\rightarrow$  LUMO + 1) + 0.48 (HOMO  $\rightarrow$  LUMO + 3). Thus, the C–CI bond dissociation would occur. Absorption bands due to CH<sub>2</sub>Cl<sub>2</sub> were

observed at 1267 and 747 cm<sup>-1</sup>, and 1263 and 746 cm<sup>-1</sup>, in the Kr and Xe matrices, respectively. The band at 853 cm<sup>-1</sup> in the Kr matrix was assigned to the v<sub>3</sub> of Kr<sub>2</sub>H<sup>+</sup> [23]. The bands at 954, 844 and 731 cm<sup>-1</sup> in the Xe matrix were assigned to the v<sub>3</sub> + 2v<sub>1</sub>, v<sub>3</sub> + v<sub>1</sub> and v<sub>3</sub> of Xe<sub>2</sub>H<sup>+</sup>, respectively [23]. For comparison, the IR difference spectrum of the photolysis products of the matrix CH<sub>2</sub>Cl<sub>2</sub>/Xe is shown in Fig. 4(c). Similar product bands were observed at 954, 844 and 731 cm<sup>-1</sup>, indicating that the photolysis product CH<sub>2</sub>Cl<sub>2</sub> of CH<sub>2</sub>ClCOCl was also further decomposed to yield Kr<sub>2</sub>H<sup>+</sup> or Xe<sub>2</sub>H<sup>+</sup>. Therefore, it is plausible that the band at 844 cm<sup>-1</sup> which was previously assigned to the v<sub>5</sub> of CHCl=C=O, is reassigned to the v<sub>3</sub> + v<sub>1</sub> of Xe<sub>2</sub>H<sup>+</sup>.



Fig. 2. Optimized equilibrium structures of the CHCI=C=O-HCI complexes calculated at the MP2/6-311++G(3df,3pd) level. Bond lengths are in Å. (a) π complex and (b) n complex



Fig. 3. Infrared difference spectra upon  $\lambda > 253.7$  nm irradiation of the matrix CH<sub>2</sub>CICOCI/Kr = 1/1000 for 360 min. (a) 330–0 min and (b) 360–330 min



Fig. 4. Infrared difference spectra upon  $\lambda > 253.7$  nm irradiation of the matrix CH<sub>2</sub>CICOCI/Xe = 1/1000. (a) 180–0 min and (b) 420–180 min. (c) Infrared difference spectra upon  $\lambda > 253.7$  nm irradiation for 180 min of the matrix CH<sub>2</sub>Cl<sub>2</sub>/Xe = 1/1000

# 3.3 CH<sub>2</sub>CICOCI/O<sub>2</sub>

To clarify the route of CHCI=C=O and CH<sub>2</sub>Cl<sub>2</sub> formation, i.e., radical or concerted mechanism, the reactive O<sub>2</sub> matrix was used. Fig. 5 shows the IR difference spectrum obtained after 300 min of the irradiation of the CH<sub>2</sub>CICOCI/O<sub>2</sub> matrix. Product bands were assigned by comparison with the spectrum observed in the photolysis of the CCl<sub>3</sub>COCI/O<sub>2</sub> matrix [13]. Due to the fact that the photolysis occurred in O<sub>2</sub> at 253.7 nm, the band of ozone was prominent at 1038 cm<sup>-1</sup> (*v*<sub>3</sub>) [22]. Other O<sub>3</sub> absorption bands were observed at 2107 (*v*<sub>1</sub> +*v*<sub>3</sub>) and 1101 cm<sup>-1</sup> (*v*<sub>1</sub>) [21,22]. The bands at 2342 and 2276 cm<sup>-1</sup> were assigned to *v*<sub>3</sub> vibrations of <sup>12</sup>CO<sub>2</sub> and <sup>13</sup>CO<sub>2</sub>, respectively. The 2037 cm<sup>-1</sup> band is attributed to CO<sub>3</sub> in complex with CI [13]. In CH<sub>2</sub>CICOCI/O<sub>2</sub>, the bands at 2157, 2152 and 835 cm<sup>-1</sup> due to CHCI=C=O were observed. Absorption bands due to CH<sub>2</sub>Cl<sub>2</sub> were observed at 1268 and 750 cm<sup>-1</sup>.

# 3.4 Reaction Mechanism

Fig. 6 shows the absorbance changes of the C=O stretching vibration band for the CHCI=C=O···HCI  $\pi$  complex and the CCI<sub>2</sub> antisymmetric stretching vibration band for CH<sub>2</sub>CI<sub>2</sub> observed in Ar, Kr and Xe. The CHCI=C=O··HCI complex showed the growth and decay profiles in Kr and Xe. Therefore, the approximate relative yield,  $\Phi_{CHCI=C=O-HCI/CH2CI2}$ , of the above mentioned absorption bands of CHCI=C=O··HCI and CH<sub>2</sub>CI<sub>2</sub> at an irradiation time of 100 min was compared using the IR intensities of 618 and 149 km mol<sup>-1</sup>, respectively, calculated at the B3LYP/aug-cc-pV(T+d)Z level. The values of  $\Phi_{CHCI=C=O-HCI/CH2CI2}$  were 18, 5.2 and 0.82 for the Ar, Kr and Xe matrices, respectively. The formation of CH<sub>2</sub>CI<sub>2</sub> was greatly enhanced in Xe probably due to the external heavy atom effect. This observation indicates that changing the matrix gas from Ar to Kr or Xe may open

the reaction paths in the triplet state, CH<sub>2</sub>ClCOCI (T<sub>1</sub>)  $\rightarrow$  CH<sub>2</sub>Cl + COCl  $\rightarrow$  CH<sub>2</sub>Cl + CO + Cl  $\rightarrow$  CH<sub>2</sub>Cl<sub>2</sub> + CO and CH<sub>2</sub>ClCOCI (T<sub>1</sub>)  $\rightarrow$  CH<sub>2</sub>ClCO + Cl  $\rightarrow$  CH<sub>2</sub>Cl<sub>2</sub> + CO.



Fig. 5. Infrared difference spectrum upon  $\lambda > 253.7$  nm irradiation of the matrix CH<sub>2</sub>CICOCI/O<sub>2</sub> = 1/1000 for 60 min



Fig. 6. Absorbance changes of CHCI=C=O and  $CH_2CI_2$  in (a) Ar, (b) Kr and (c) Xe

As shown in Fig. 5, the ketene species were produced in  $O_2$ , which was indicative of their formation through a concerted mechanism in the ground state. A small amount of  $CH_2Cl_2$  was also observed in  $O_2$ . The energy diagram for the formation of the ketene species and  $CH_2Cl_2$  in the ground state is shown in Fig. 7. The transition state TS3 for the formation of  $CH_2Cl_2$  and CO lies in much higher compared with the TS2 for the formation of CHCI=C=O—HCI. This agreed with the observed predominant CHCI=C=O formation in Ar.

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# Fig. 7. Energy diagram for the reactions of CH<sub>2</sub>CICOCI in the ground state calculated at the B3LYP/aug-cc-pV(T+d)Z level

## 4. CONCLUSION

UV light photolysis of CH<sub>2</sub>CICOCI was investigated in cryogenic Ar, Kr, Xe and O<sub>2</sub> matrices. In CH<sub>2</sub>CICOCI/Ar, photoisomerization from *anti*- to *gauche*-CH<sub>2</sub>CICOCI was observed at an early stage of the irradiation. In all matrices, CHCI=C=O, HCI, CH<sub>2</sub>Cl<sub>2</sub> and CO were produced, with the relative yields depending on the used matrix gas The relative yield of CH<sub>2</sub>Cl<sub>2</sub> formation was greatly enhanced in Xe because of the external heavy atom effect. The results show that in Ar, the reactions predominantly proceed through concerted mechanisms.

# **COMPETING INTERESTS**

Authors have declared that no competing interests exist.

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